



Meenakshi Narain Brown University



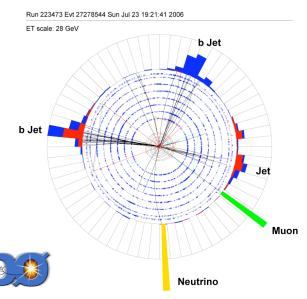
LHC@BNL workshop June 19, 2009

top at Fermilab

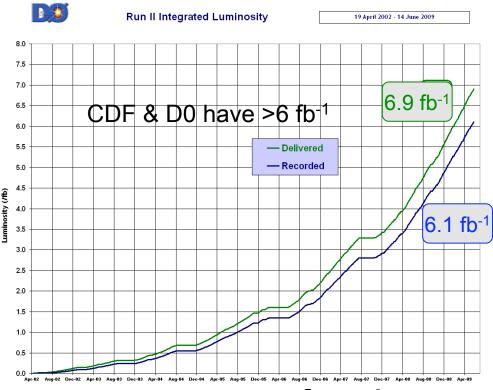
- 14 years ago...
 - ...we observed a few handfuls of top quarks.
- Recently...

Observed EWK production process: single top quark

DØ Experiment Event DisplaySingle Top Quark Candidate Event, 2.3 fb⁻¹ Analysis



today...



...we are performing detailed studies of 1000s of top decays

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outline

- strong production
 - cross section

branching fractions

antiproton beam

W

electron

- new physics?
 - FCNC decays
 - tt resonances
 - tb resonances

mass

muon

proton beam

neutrino

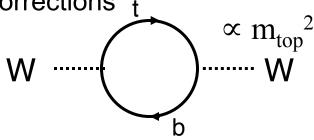
couplings

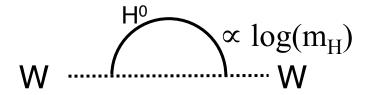
electroweak production

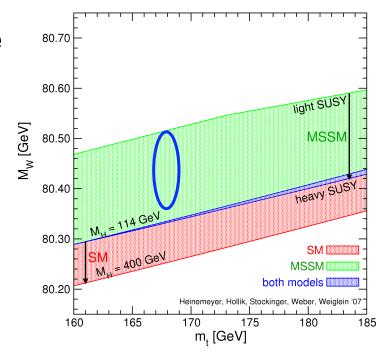
 $-|V_{tb}|$

why is the top quark important?

- most massive elementary particle
 - dominant contributor to radiative corrections t



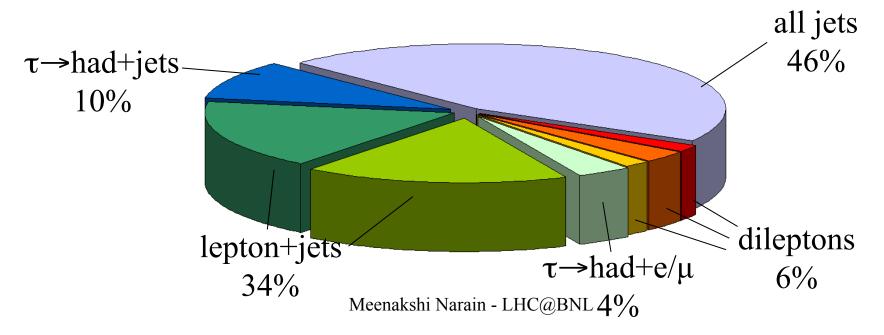


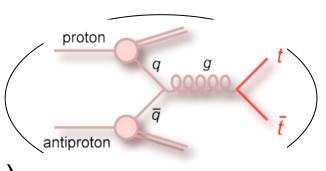


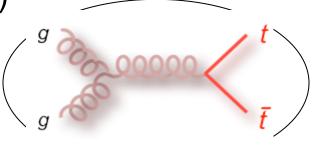
- how is its mass generated?
 - topcolor?
- does it couple to new physics?
 - massive G, heavy Z', H+, ...

top-antitop production

- strong interaction
 - \rightarrow top-antitop pairs ($\sigma = 7.6 \pm 0.6 \text{ pb}$)
- final state signatures for top-antitop pairs
- $t\rightarrow Wb$ with $B \approx 100\%$
 - →tagging b-jets important
 - W→qq with $B \approx 67$ %; W→ ℓ_V with $B \approx 11$ %
 - τ → evv/μvv with $B \approx 17 \%$







Tevatron phases

- Run la (20 pb⁻¹ handful of events)
 - e-mu is the golden channel to find top
 - if top is really massive will we be able to see it?
- Run Ib (160 pb⁻¹ <100 events)
 - top quark is really massive
 - need to use the hadronic signatures
 - I+jets is the golden channel
 - measure cross section and mass
- Run IIa (1 fb⁻¹ 100s of events)
 - advanced analysis methods
 - more properties measured
 - precision measurements
 - combination of all measurements of an observable
- Run IIb (8? fb⁻¹ 1000s events)
 - beginning to get systematically limited
 - consistency of measurements of different observables
 - look at the whole picture

Lessons learned

- Be prepared to abandon your premise
 - To find the top quark it was necessary to re-optimize the search for much higher mass using different channels and features in the event than initially imagined.
 - To optimize sensitivity and resolution it was necessary to invent new analysis techniques.
 - b-tagging
 - Mass measurement in dilepton channel
 - Matrix element/event-by-event likelihood techniques
 - In situ calibration of jets with W->qq
 - Neural networks, decision trees, etc.

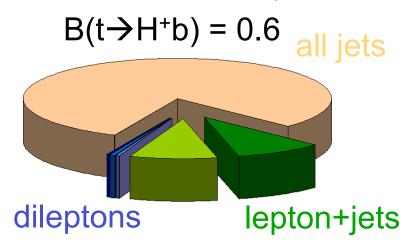
Lessons learned

- Develop a simulation that mimics detector performance
 - Crucial to extrapolate from e.g. background to signal regions
 - Study systematic effects
 - Realize what you cannot simulate well (jet multiplicity, fake rates)
 - Often limited by rate at which events could be simulated
 - need for fast MC simulation tuned to data
- Control samples to verify distributions, estimate backgrounds, efficiencies for triggers, lepton-id etc.

Selected Results from Dzero

why measure the ttbar cross section?

- cross section analysis
 - basic understanding of signal and background necessary for further study
 - consistency between channels
 - decay branching fractions
 - are there non-standard decays?



ttbar cross section in I+jets channel

- extract top fraction using event topology
 - angles, momentum sums, and event shape variables
 - dominated by statistical uncertainties
- count number of events with at least one b-tagged jet
 - smaller statistical uncertainty
 - large systematic uncertainty from jet energy calibration and b-tagging

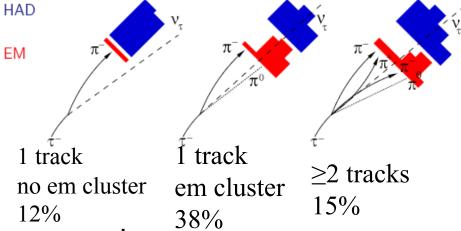
 $D0 (0.9 \text{ fb}^{-1})$ $D0 (0.9 \text{ fb}^{-1})$ σ =8.1±0.5(stat)±0.7(syst)±0.5(lum) pb σ =6.6±0.8(stat)±0.4(syst)±0.4(lum) pb D0 Run II 0.9 fb⁻¹ D0 Run II 0.9 fb⁻¹ Number of Events 000 000 000 000 000 000 000 data (1 b-tag) data (≥4 jets) tt signal tt signal bkg with leptons bkg with leptons bkg w/o leptons bkg w/o leptons 2 Jet Multiplicity **Likelihood Discriminant**

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ttbar cross section in τ channels

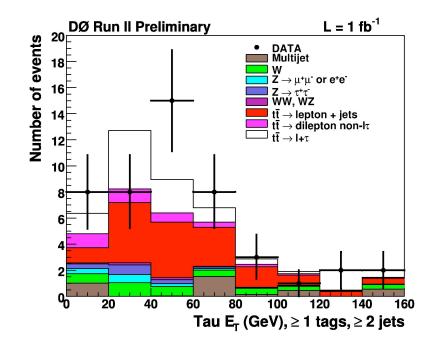


- interesting because of t→H⁺b, H⁺→τν
- 3 types of hadronic τ decays



- require
 - -1τ , 1 e/μ, ≥2 jets, missing p_T
 - 1 jet is b-tagged

 neural networks distinguish τ decays from background



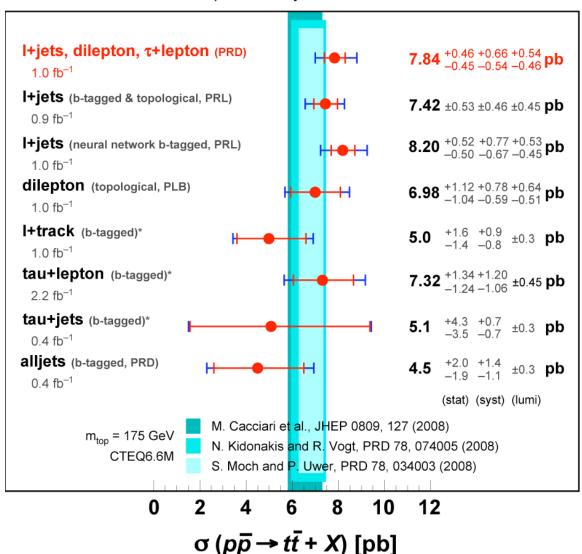
$$\sigma = 7.32^{+1.3}_{-1.2} (\text{stat})^{+1.2}_{-1.1} (\text{syst}) \pm 0.4 (\text{lum}) \text{ pb } (2.2 \text{ fb}^{-1})$$

$$\sigma B(\text{tt} \rightarrow \text{l}\tau) = 0.19 \pm 0.08 (\text{stat}) \pm 0.07 (\text{syst}) \pm 0.01 (\text{lum}) \text{ pb } (1 \text{ fb}^{-1})$$

ttbar cross section summary

DØ Run II * = preliminary

May 2009



Source	$\Delta \sigma_{t\bar{t}}$ (pb)
Statistical only	+0.47 -0.46
Lepton identification	+0.15 -0.14
Tau identification	+0.02 -0.02
Jet identification	+0.11 -0.11
Jet corrections	+0.19 -0.16
Tau energy scale	+0.02 -0.02
Trigger modeling	+0.11 -0.07
b jet identification	+0.34 -0.32
Signal modeling	+0.17 -0.15
Background estimation	+0.14 -0.14
Multijet background	+0.12 -0.12
Luminosity	+0.56 -0.48
Other	$+0.15 \ -0.14$
Total systematic uncertainty	+0.78 -0.69

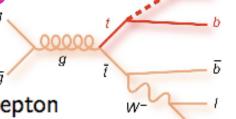
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non standard decay modes

New particles in final state may alter σ(tt):

$$\sigma_{t\bar{t}}^{\mathrm{Ch}} = \sigma_{t\bar{t}} \cdot \frac{B^{\mathrm{BSM}}(t\bar{t} \to \mathrm{Ch})}{B^{\mathrm{SM}}(t\bar{t} \to \mathrm{Ch})}$$

 $\mathrm{Ch} \!=\! \ell \!+\! \mathrm{jets}$, Dilepton, $\tau \!+\! \mathrm{lepton}$

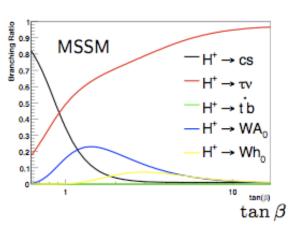


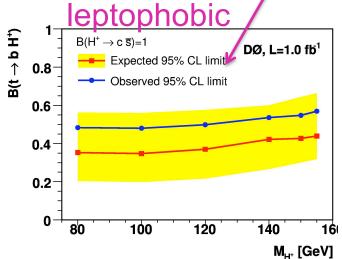
Check cross section ratios

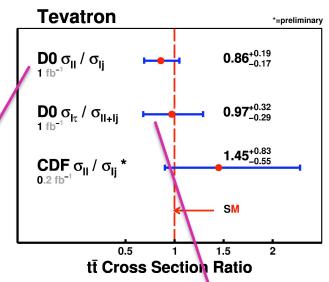
$$\sigma_{t\bar{t}}^{\ell+\mathrm{jets}}/\sigma_{t\bar{t}}^{\mathrm{Dilepton}}/\sigma_{t\bar{t}}^{\tau+\ell}$$

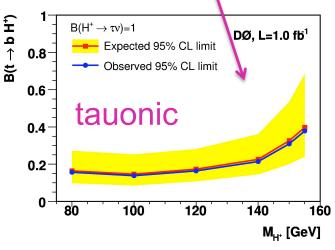
- Many systematic uncertainties are canceled in the ratios
- Within MSSM depending on tanβ

 $H^{\pm} \rightarrow cs$ or $\tau \nu$ may dominate.





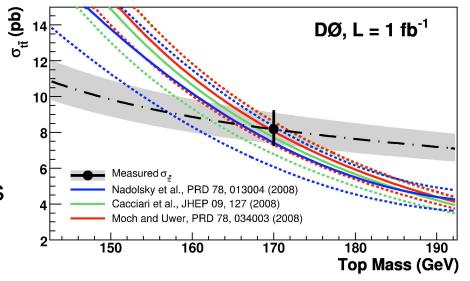




implications of cross section



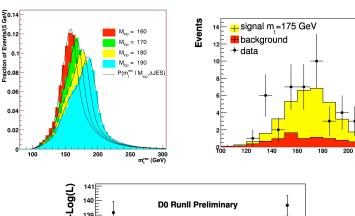
- compare with theory to obtain mass measurement
 - Less sensitive to the non-perturbative QCD effects
- Construct likelihood with measurements and define theory likelihoods according to PDF and scale uncertainties as in refs 1-4.
- Determine joint theory and expt'al likelihood, integrate over the cross section & get 68% CL.
- Nadolsky et. al., Phys. Rev. D 78 013004 (2008); W. Beenakker et. al., Phys. Rev. D 40, 54 (1989).
- 2. Cacciari et al., jEHP 09, 127 (2008).
- 3. Moch & Uwer, Phys. Rev. D 78 034003 (2008).
- 4. Kidonakis & Vogt, Phys. Rev. D 78 074005 (2008).

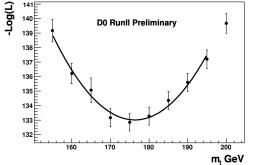


Theoretical computation	$m_{\rm t}~({\rm GeV})$
NLO [1]	$165.5_{-5.9}^{+6.1} 167.5_{-5.6}^{+5.8} 169.1_{-5.2}^{+5.9}$
NLO+NLL [2]	$167.5^{+5.8}_{-5.6}$
approximate NNLO [3]	$169.1^{+5.9}_{-5.2}$
approximate NNLO [4]	$168.2^{+5.9}_{-5.4}$

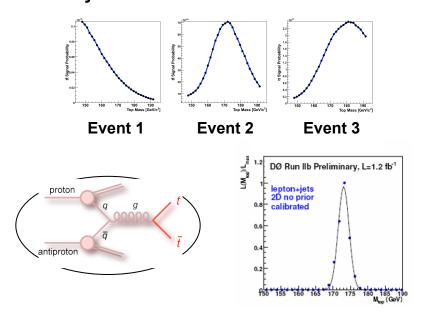
top mass measurement

- template fits
 - mass estimator (eg best m_t from kinematic fitter)
 - fit probability density functions from simulated tt events and background to data





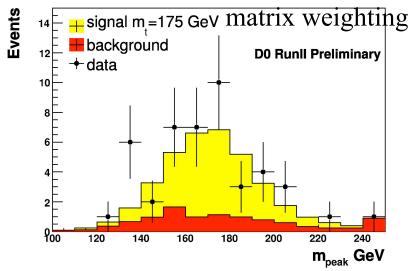
- event-by-event likelihood
 - for each event determine likelihood as a function of m_t (eg by integrating over LO matrix element)
 - extract mass from peak of joint likelihood



 $\mathbf{m}_{\mathrm{peak}}$ GeV

dilepton channel

- D0 (1 fb⁻¹)
- matrix weighting and neutrino weighting techniques
- compute weight curve as a function of top mass for each event
- template fit to mass distribution
- Combined measurement:

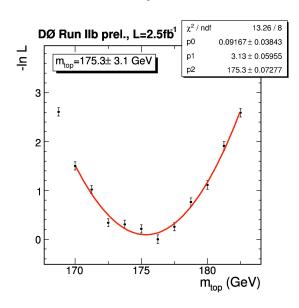


174.7±4.4(stat)±2.0(syst) GeV

Source of uncertainty	νWT_h	MWT
	[GeV]	[GeV]
b fragmentation	0.4	0.4
Underlying events modeling	0.3	0.5
Extra jets modeling	0.1	0.3
Event generator	0.6	0.5
PDF variation	0.2	0.5
Background template shape	0.4	0.3
Jet energy scale	1.6	1.2
b/light response ratio	0.3	0.6
Sample dependent JES	0.4	0.1
Jet resolution	0.1	0.2
Muon/track resolution	0.1	0.2
Electron resolution	0.1	0.2
Jet identification	0.4	0.5
MC corrections	0.2	0.2
Background yield	0.0	0.1
Signal shape modeling	0.8	0.8
MC calibration	0.1	0.1
Total systematic uncertainty	2.1	2.0

dilepton channel

- D0 (3.6 fb⁻¹)
 - compute event weight using LO matrix element
 - Use electron-muon events
 - Clean sample, little background



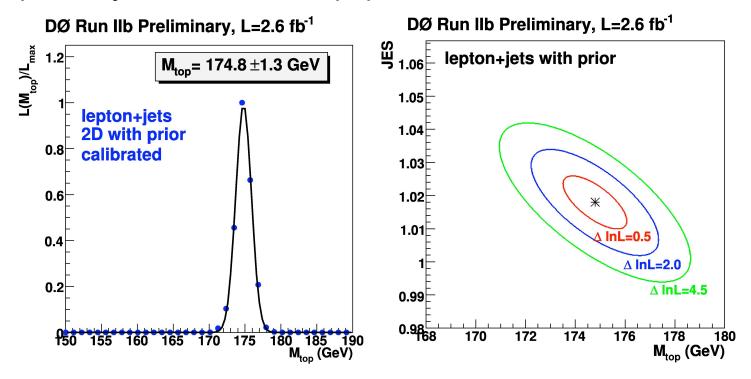
Uncertainty	$e\mu$ Run IIb [GeV]
JES up	-1.5
JES down	+1.8
b quark JES	+1.4
jet resolution up	-0.7
jet resolution down	+0.7
jssr shifting	+0.1
muon smearing up	-0.0
muon smearing down	+0.3
b quark fragmentation	± 0.3
PDF uncertainty up	-0.2
PDF uncertainty down	+0.1
fit uncertainty	± 0.4
signal modeling	± 0.4
background fraction up	-0.1
background fraction down	+0.2
Total	$^{+2.5}_{-1.8}$

174.7±2.9(stat)±2.4(syst) GeV

best precision

lepton+jets

- matrix element analysis (3.6 fb⁻¹)
 - integrate over LO matrix element to get likelihood for event as a function of top quark mass
 - in situ jet energy calibration using W→qq decay
 - peak of joint likelihood = top quark mass



D0: 173.7±0.8(stat)±1.6(syst⊕jes) GeV Meenakshi Narain - LHC@BNL

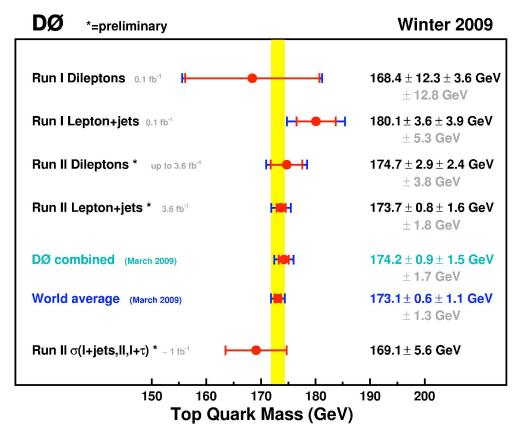
Lepton+jets systematics:

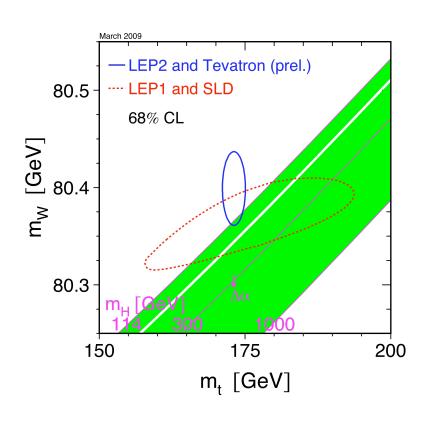
D0

Source	Uncertainty (Ge	V)
Higher Order Effects	±(0.25
ISR/FSR	±(0.26
Hadronization and UE	±(0.58
Color Reconnection	±(0.50
PDF uncertainty	±(0.24
Residual JES uncertainty	±(0.21
Relative b /light response	±(0.81
Sample-dependent JES	±(0.56
Jet ID efficiency	±(0.26
Jet energy resolution	±(0.32
Plus a few smaller sys <0.2	ı	
Total	±1	1.44

Combination (as of winter 09)

D0 winter '09 $m_{top} = 174.2 \pm 0.9(stat) \pm 1.5(syst)$ GeV





 $\delta m/m < 1\%$

Run II goal: δm ≈ 1 GeV

http://tevewwg.fnal.gov/top/

http://lepewwg.web.cern.ch/LEPEWWG/plots/winter2009/

theory issues

- what mass are we measuring?
 - Pole mass? (direct msm't calibrate to MC).
 - understanding required for consistency checks and M_H prediction
 - EW precision fits use MS mass
- are we missing any important effects?
 - Color reconnection
- for the LHC, the complementary approach from measurements using cross section may have the potential to eventually get to similar level of systematic uncertainties?

analysis trends

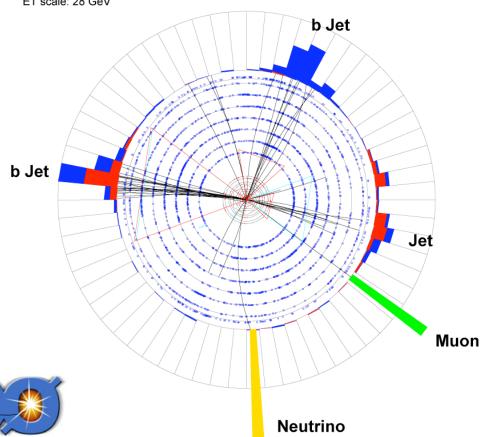
- multivariate analysis methods
 - needed to get optimal sensitivity
 - require good understanding of detector and simulation

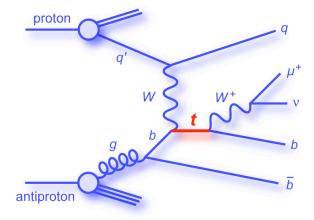
observation of single top production

DØ Experiment Event Display

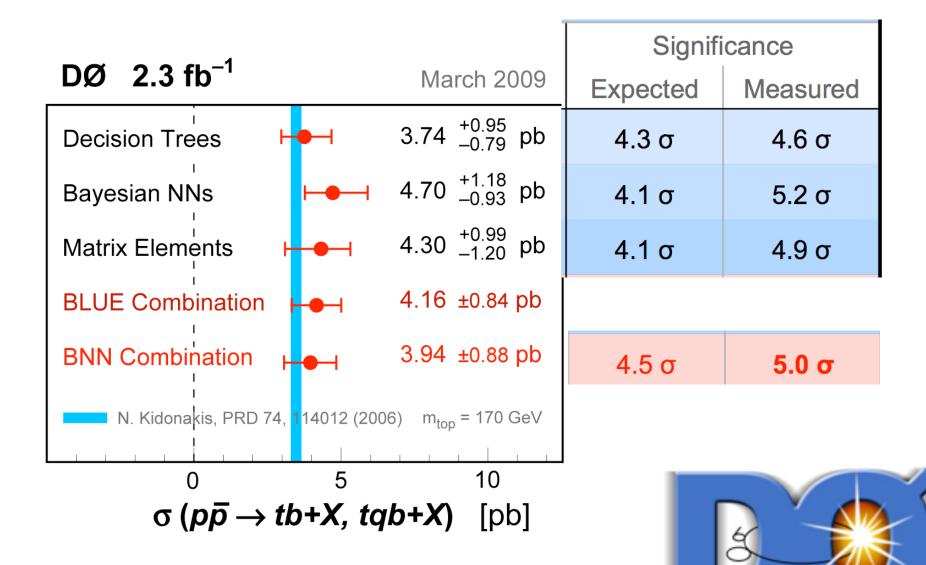
Single Top Quark Candidate Event, 2.3 fb⁻¹ Analysis

Run 223473 Evt 27278544 Sun Jul 23 19:21:41 2006 ET scale: 28 GeV b Jet

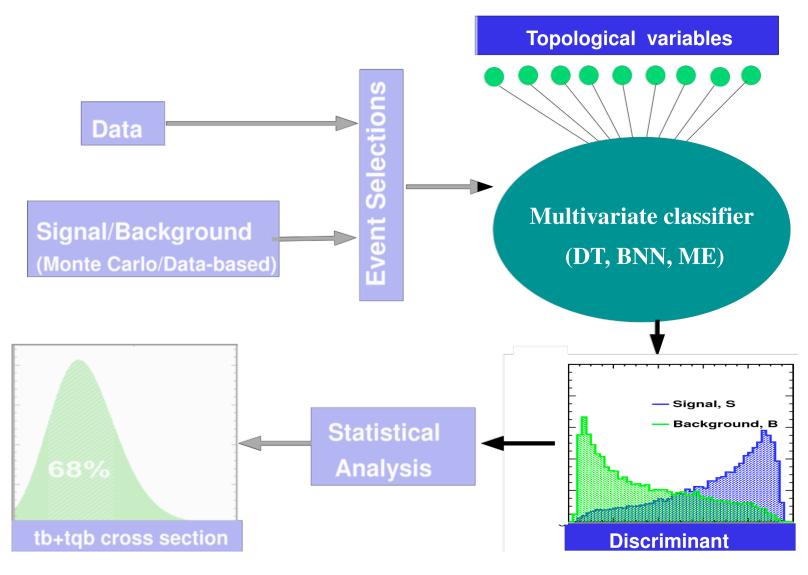




cross section summary



analysis flow chart



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event counts - final selection:

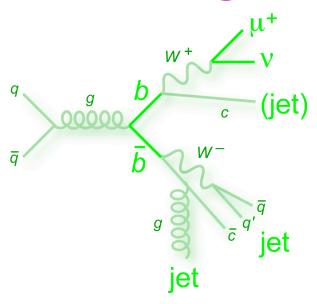
- expected signal
- backgrounds

observed

Event Yields in 2.3 fb ⁻¹ of DØ Data	
e,μ, 2,3,4-jets, 1,2-tags combined	
tb + tqb	223 ± 30
<i>W</i> +jets	2,647 ± 241
Z+jets, dibosons	340 ± 61
<i>tt</i> pairs	1,142 ± 168
Multijets	300 ± 52
Total prediction	4,652 ± 352
Data	4,519

modeling of backgrounds:

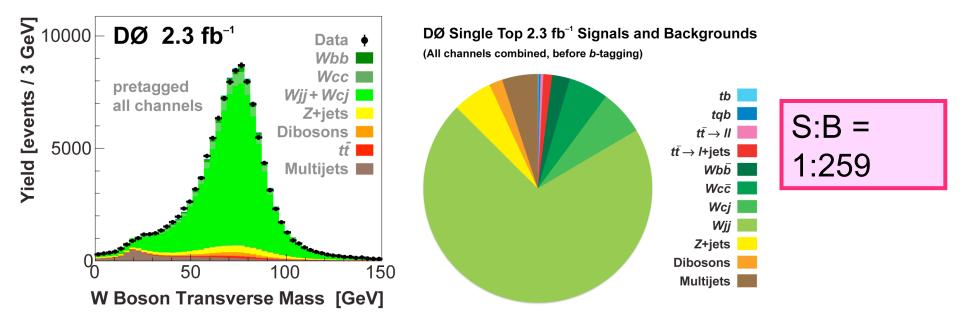
W+jets:



- modeled using ALPGEN
 - PYTHIA for parton hadronization
 - MLM parton-jet matching avoids double-counting final states
- η(jets), Δφ(jet1,jet2), Δη(jet1,jet2) corrected to match data

1. bkg normalization pre b-tagging

dominant background: W+jets

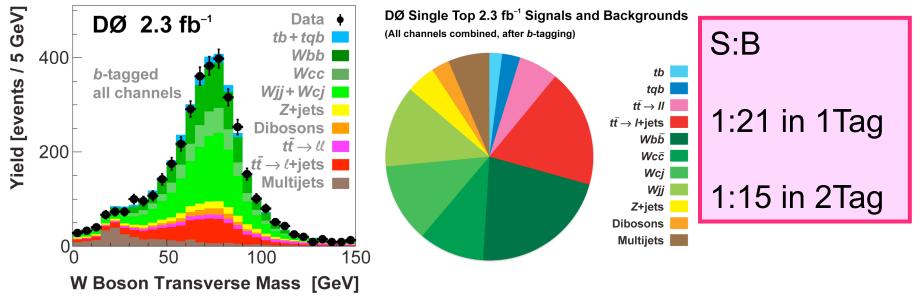


 Overall normalization for Wjets/mis-id determined by using iterative template fits to data using three sensitive variables: p_T(I), M_T(W) and missing E_T

$$N_{\mathrm{pretag}}^{\mathrm{data}} - N_{\mathrm{bkgd}}^{\mathrm{MC}} = S_{\mathrm{W+jets}} N_{\mathrm{W+jets}}^{\mathrm{MC}} + S_{\mathrm{multijet}} N_{\mathrm{multijet}}^{\mathrm{data}}$$
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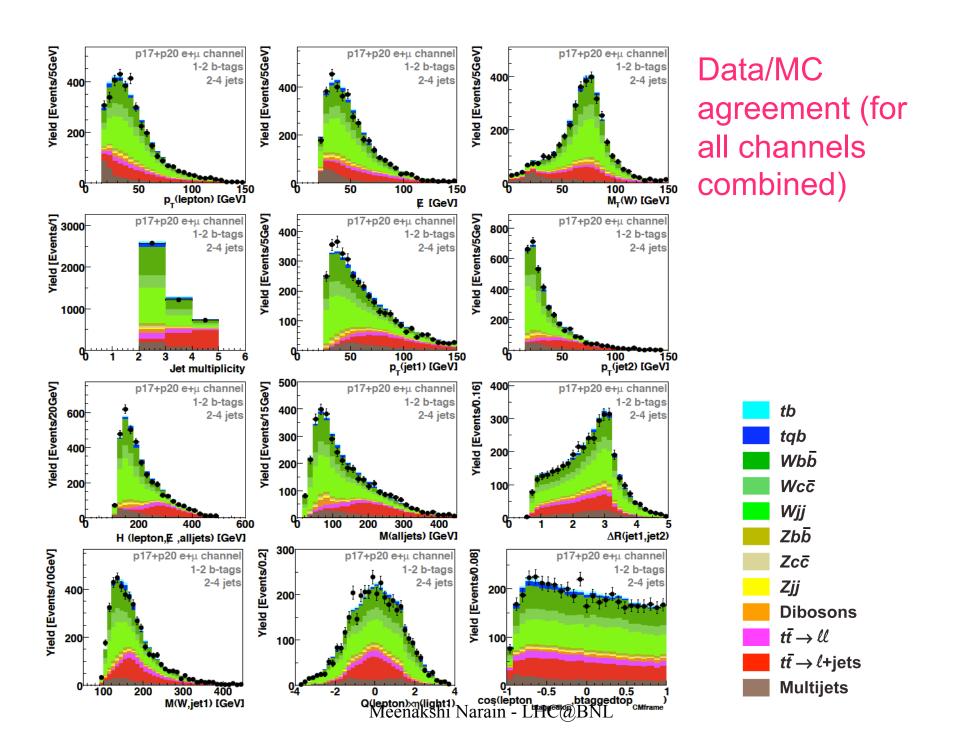
2. bkg norm. post b-tagging

 W+HF (Wbb, Wcc, Wcj), top pair backgrounds are dominant



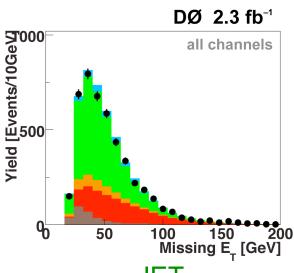
- W+heavy flavor correction factors
 - normalized to theory (use MCFM @ NLO)
 - 1.47 (Wbb, Wcc), 1.38 (Wcj)
 - additional empirical correction derived from two-jet data and simulation: includes zero-tag events
 - 0.95 ± 0.13 (Wbb, Wcc)

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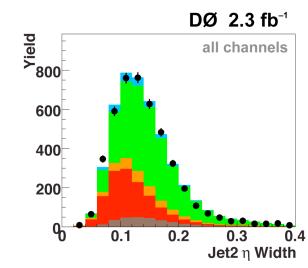


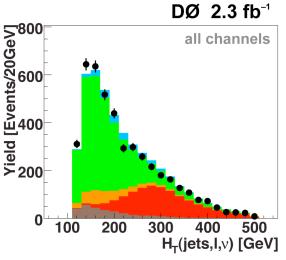
discriminating variable categories



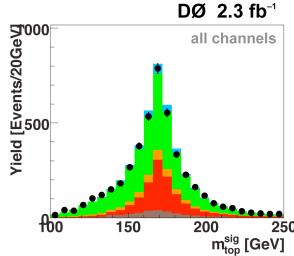


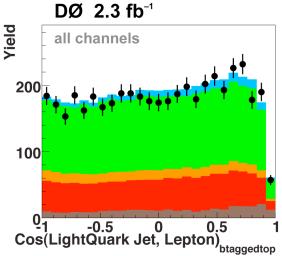
JET RECONSTRUCTION



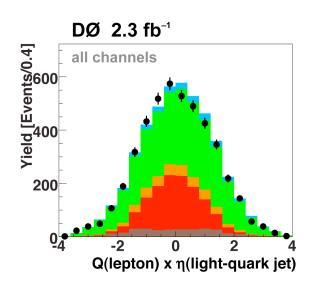


TOP QUARK RECONSTRUCTION



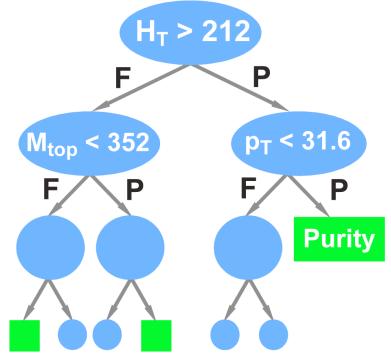


ANGULAR CORRELATIONS



boosted decision trees

- decision trees
 - idea: recover events that fail a cut
 - successively find cuts with best separation between signal and background
 - repeat recursively on each branch
 - stop when no further improvement or when too few events are left
 - terminal node is called a "leaf"
 - decision tree output = leaf purity



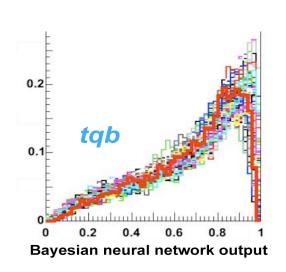
boosted decision trees

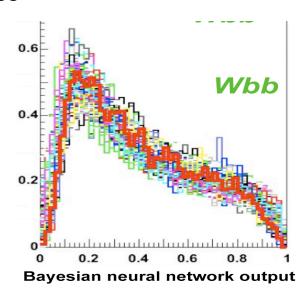
Best Variables to Separate Single Top from W+Jets	
DØ 2.3 ft	o ^{−1} Analysis
Object kinematics	# _⊤
	ρ _τ (jet2)
	p_T^{rel} (jet1,tag- μ)
	E(light1)
Event kinematics	M(jet1,jet2)
	$M_{7}(W)$
	H_T (lepton, $\not\equiv_T$,jet1,jet2)
	H ₇ (jet1,jet2)
	H_T (lepton, $\not\equiv_T$)
Jet reconstruction	Width _o (jet2)
	Width _η (jet2)
Top quark reconstruction	M _{top} (W,tag1)
	$\Delta M_{\mathrm{top}}^{\mathrm{min}}$
	M _{top} (W,tag1,S2)
Angular correlations	cos(light1,lepton) _{btaggedtop}
	Δφ(lepton, ∉ _τ)
	Q(lepton) x η(light1)

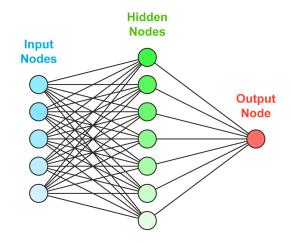
Best Variables to Separate Single Top from Top Pairs	
DØ 2.3 ft	o ^{–1} Analysis
Object kinematics	pT(notbest2)
	<i>pT</i> (jet4)
	pT(light2)
Event kinematics	M(alljets-tag1)
	Centrality(alljets)
	M(alljets-best1)
	H_{τ} (alljets–tag1)
	$H_{\tau}(lepton, \not\!\!\!\!/_{\tau}, alljets)$
	M(alljets)
Jet reconstruction	Width _η (jet4)
	Width _o (jet4)
	Width _{ϕ} (jet2)
Angular correlations	$cos(lepton_{btaggedtop}, btaggedtop_{CMframe})$
	Q(lepton) x η(light1)
	ΔR (jet1,jet2)

bayesian neural networks

- Neural networks are nonlinear functions
 - defined by weights associated with each node
 - weights are determined by training on signal and background samples
- Bayesian neural networks improve on this technique
 - average over many networks weighted by the probability of each network given the training samples
 - Less prone to over-training
 - Network structure is less important can use larger numbers of variables and hidden nodes







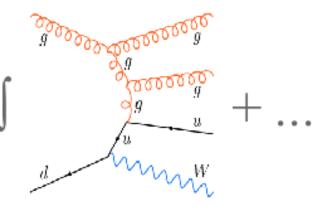
bayesian neural networks

- list of variables
 - example from one channel.

Rank	Variable	
1	M(jet1,jet2)	
2	$M_T(W)$	
3	$M_{top}(W, \text{tag1})$	
4	$\Delta \ { m M}_{top}^{min}$	
5	$H_T(\text{lepton}\cancel{L}_T, \text{jet1, jet2})$	
6	$M_{top}(W, \text{tag1}, S2)$	
7	$ \not\!\!E_T$	
8	$Q(lepton) \times \eta(light1)$	
9	$\cos(\mathrm{lepton_{btaggedtop}}, \mathrm{btaggedtop_{CMframe}})$	
10	$\cos(\text{tag1}, \text{lepton})_{\text{btaggedtop}}$	
11	$p_T(\text{jet }1)$	
12	$Width_{\eta}(jet2)$	
13	$\Delta\phi(\mathrm{lepton} E_T)$	
14	$\mathrm{Width}_\phi(\mathrm{jet}2)$	
15	$p_T(\text{jet2})$	
16	$Q(lepton) \times \eta(best 1)$	
17	E(jet2)	
18	$p_T(\text{best1})$	
19	$p_T^{rel}({ m jet1},\mu)$	
20	$\cos(\mathrm{light1,lepton})_{\mathrm{btaggedtop}}$	
21	$\cos(\text{lepton}, Q(\text{lepton}) \times z)_{\text{besttop}}$	

matrix elements

- method pioneered by DØ for top quark mass measurement
- use 4-vectors of reconstructed leptons and jets
- use matrix elements of main signal and background processes
- compute a discriminant



$$D_s(\vec{x}) = P(S|\vec{x}) = \frac{P_{Signal}(\vec{x})}{P_{Signal}(\vec{x}) + P_{Background}(\vec{x})}$$

• uenne r_{signal} as a normanzeu umerenti cross section:

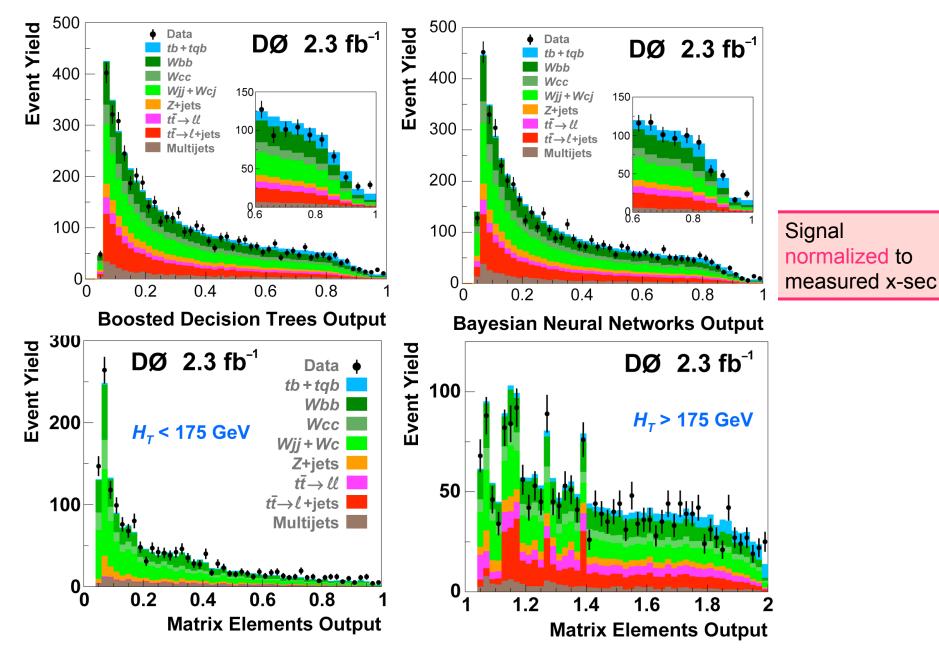
$$P_{Signal}(\vec{x}) = \frac{1}{\sigma_S} d\sigma_S(\vec{x}) \quad \sigma_S = \int d\sigma_S(\vec{x})$$

2-jets and 3-jets channels only

Matrix Elements used to Separate Single Top Signal from Background						
DØ 2.3 fb ⁻¹						
2 Jets		3 Jets				
t b	$u\overline{d} o t\overline{b}$	tbg	$u\overline{d} ightarrow t\overline{b}g$			
tq	$egin{aligned} ub & ightarrow td \ dar{b} & ightarrow tar{u} \end{aligned}$	tqg	$egin{aligned} ub & ightarrow tdg \ dar{b} & ightarrow tar{u}g \end{aligned}$			
		tq b	$egin{aligned} ug ightarrow tdar{b} \ ar{d}g ightarrow tar{u}ar{b} \end{aligned}$			
Wbb	$u\bar{d} \rightarrow Wb\bar{b}$	Wbbg	$u\overline{d} o Wb\overline{b}g$			
Wcg	ar s g o W ar c g					
Wgg	$u \overline{d} o W g g$	Wūgg	$\overline{u}g ightarrow W\overline{u}gg$			
WW	$q \overline{q} o WW$					
WZ	$q\overline{q} o WZ$					
<i>ggg</i>	gg o ggg					
tt	$q\overline{q} ightarrow t\overline{t} ightarrow \ell^+ v b \ell^- v \overline{b}$					
tt	$q \overline{q} ightarrow t \overline{t} ightarrow \ell^+ v b \overline{u} d \overline{b}$	tt	$q \overline{q} ightarrow t \overline{t} ightarrow \ell^+ v b \overline{u} d \overline{b}$			

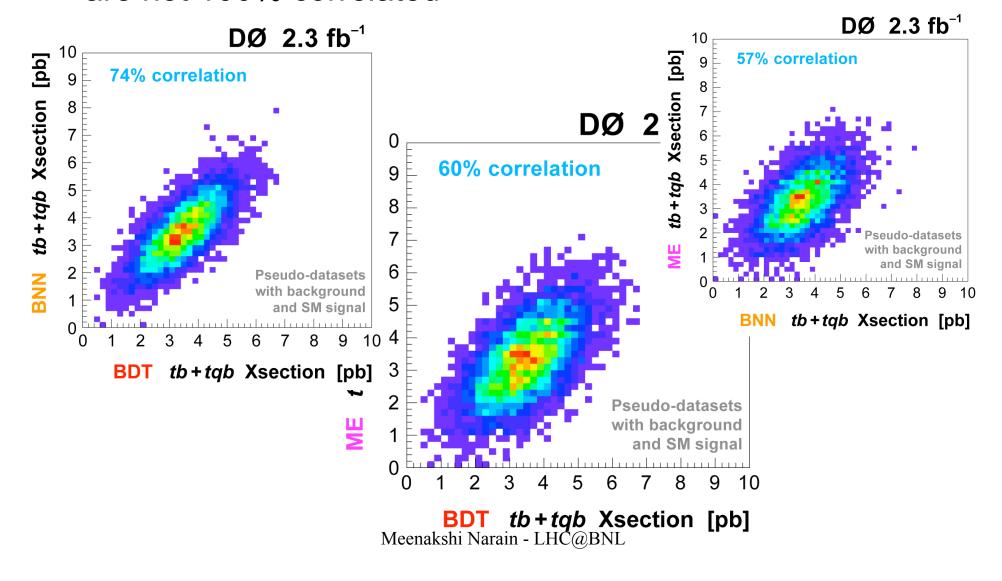
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final discriminant for the 3 methods



correlations between methods

 Even though all MVA analyses use the same data, they are not 100% correlated



Systematic uncertainties

Systematic Uncertainties

Ranked from Largest to Smallest Effect on Single Top Cross Section

DØ 2.3 fb⁻¹

Larger terms				
<i>b</i> -ID tag-rate functions (includes shape variations)	(2.1–7.0)% (1-tag) (9.0–11.4)% (2-tags)			
Jet energy scale (includes shape variations)	(1.1–13.1)% (signal) (0.1–2.1)% (bkgd)			
W+jets heavy-flavor correction	13.7%			
Integrated luminosity	6.1%			
Jet energy resolution	4.0%			
Initial- and final-state radiation	(0.6–12.6)%			
b-jet fragmentation	2.0%			
$tar{t}$ pairs theory cross section	12.7%			
Lepton identification	2.5%			
Wbb/Wcc correction ratio	5%			
Primary vertex selection	1.4%			

Systematic Uncertainties

Ranked from Largest to Smallest Effect on Single Top Cross Section

DØ 2.3 fb⁻¹

Smaller terms				
Monte Carlo statistics	(0.5–16.0)%			
Jet fragmentation	(0.7–4.0)%			
Branching fractions	1.5%			
Z+jets heavy-flavor correction	13.7%			
Jet reconstruction and identification	1.0%			
Instantaneous luminosity correction	1.0%			
Parton distribution functions (signal)	3.0%			
Z+jets theory cross sections	5.8%			
W+jets and multijets normalization to data	(1.8–3.9)% (W+jets) (30–54)% (multijets)			
Diboson theory cross sections	5.8%			
Alpgen W+jets shape corrections	shape only			
Trigger	5%			

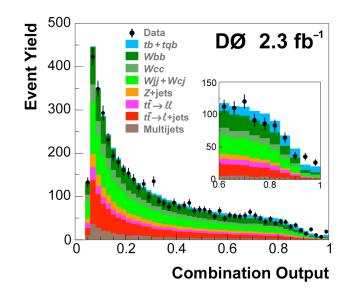


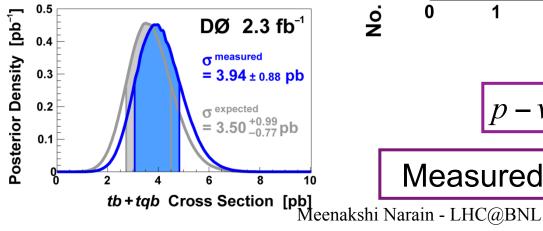
Combined results

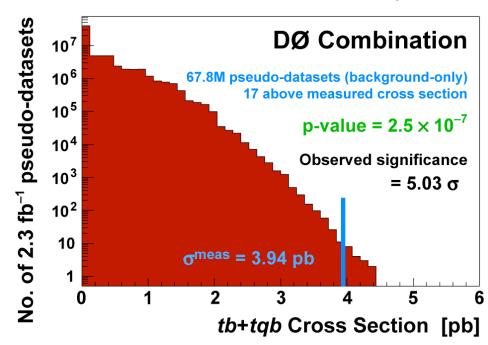


$$\sigma(p\overline{p} \rightarrow tb + X, tqb + X) = 3.94 \pm 0.88 \text{ pb}$$

 $(m_t=170GeV)$







$$p - value = 2.5 \times 10^{-7}$$

Measured Significance = 5.03σ

Lessons learned

- Lot of detailed work goes into making sure the data and MC agree.
- Control samples to verify distributions, estimate backgrounds, efficiencies for triggers, lepton-id etc.
- Validation of the inputs to the MVAs and output discriminants are well modeled.

CKM matrix element Vtb

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = V_{CKM} \begin{pmatrix} d \\ s \\ b \end{pmatrix} \qquad V_{CKM} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \qquad t \qquad V_{tb}$$

- Weak interaction eigenstates and mass eigenstates are not the same: there is mixing between quarks, described by CKM matrix
- general form of Wtb vertex:

$$\Gamma^{\mu}_{Wtb} = -\frac{g}{\sqrt{2}} \ V_{tb} \ \left\{ \gamma^{\mu} \left[f_1^L P_L + f_1^R P_R \right] - \frac{i\sigma^{\mu\nu}}{M_W} \left(p_t - p_b \right)_{\nu} \left[f_2^L P_L + f_2^R P_R \right] \right\}$$

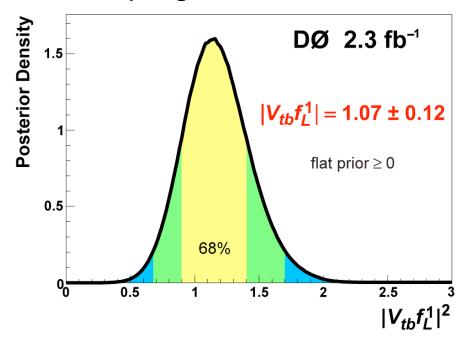
- assume
 - sm top quark decay : $|V_{td}|^2 + |V_{ts}|^2 \ll |V_{tb}|^2$
 - pure V-A : $f_1^R = 0$
 - CP conservation : $f_2^L = f_2^R = 0$
- do not assume
 - three quark families
 - CKM matrix unitarity
 - (unlike for measurements using tt decays)
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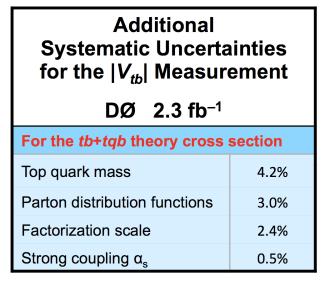
measurement of |V_{tb}|

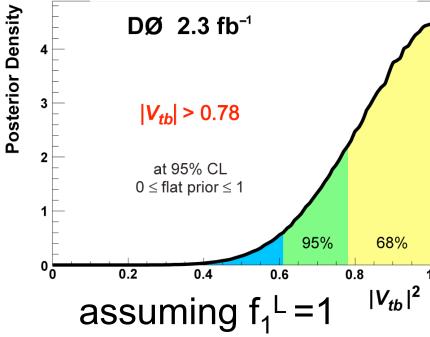
 Use the measurement of the single top cross section to make a direct measurement of |V_{tb}|:

$$\sigma(\text{tb, tqb}) \propto \left| V_{\text{tb}} f_1^L \right|^2$$

- Calculate a posterior in |V_{tb}f₁^L|²
- Measure the strength of the V–A coupling, which can be > 1







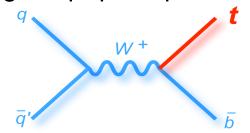
top quark coupling

- if top plays a special role in ewk symmetry breaking its couplings to W bosons may differ from predictions
- modifications to top quark interactions, in particular with weak gauge bosons, could yield the first signs of new physics
- most general CP-conserving W-t-b vertex involves four couplings

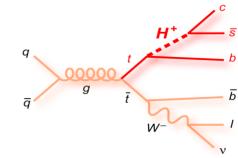
$$L_{tWb} = \frac{g}{\sqrt{2}} W_{\mu}^{-} \bar{b} \gamma^{\mu} \Big(f_{1}^{L} P_{L} + f_{1}^{R} P_{R} \Big) t - \frac{g}{\sqrt{2} M_{W}} \partial_{\nu} W_{\mu}^{-} \bar{b} \sigma^{\mu\nu} \Big(f_{2}^{L} P_{L} + f_{2}^{R} P_{R} \Big) t$$
 where, in the SM $f_{1}^{L} \approx 1$, $f_{2}^{L} = f_{1}^{R} = f_{2}^{R} = 0$ +h.c.

probing tWb vertex:

Anomalous couplings in single top quark production and decay



W helicity In top pair decays



Both measurements can be combined to fully specify the tbW vertex

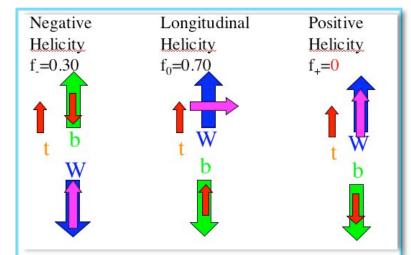
W boson helicity from t→Wb decays in top pair production

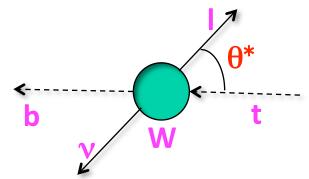
- sm predicts V-A coupling at Wtb
- ⇒ helicity of W boson

$$f_0 = 0.7$$
, $f_{-} = 0.3$, $f_{+} = 0.0$

(longitudinal, left-handed, right-handed)

- a different Lorentz structure of the t → Wb interaction would alter the fractions of W bosons produced in each polarization state.
- model-independent measurement based on reconstruction of cosθ* distribution angle between lepton and top in W rest frame
- distribution of cos θ* depends on the W boson helicity fractions

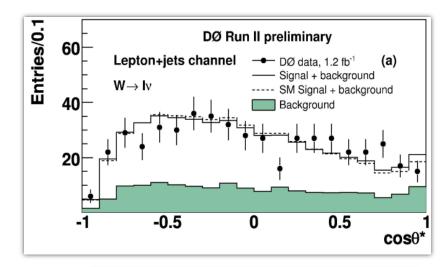




of $cos\theta^*$ (angle between the momenta of the down-type fermion and the top quark in the W boson rest frame for each top quark decay.)

top quark coupling

- Use a maximum likelihood fit, for the data to be consistent with the sum of signal and background in the cosθ* distribution
- The fit parameters are the W helicity fractions f₀ and f₊



 A model-independent measurement of the helicity of W bosons

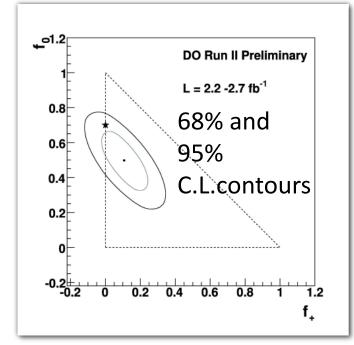
$$f_0 = 0.490 \pm 0.106 \text{ (stat.)} \pm 0.085 \text{ (syst.)}$$

 $f_+ = 0.110 \pm 0.059 \text{ (stat.)} \pm 0.052 \text{ (syst.)}$

if f₀ constrained to the standard model value

$$f_{+} = 0.019 \pm 0.031 \text{ (stat.)} \pm 0.047 \text{ (syst.)}$$

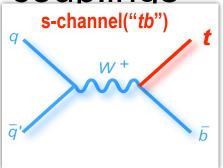
This is <u>the</u> most precise such measurement



anomalous couplings in single top production

Left & Right handed Vector and Tensor couplings

$$L_{tWb} = \frac{g}{\sqrt{2}} W_{\mu}^{-} \bar{b} \gamma^{\mu} \left(f_{1}^{L} P_{L} + f_{1}^{R} P_{R} \right) t - \frac{g}{\sqrt{2} M_{W}} \partial_{\nu} W_{\mu}^{-} \bar{b} \sigma^{\mu\nu} \left(f_{2}^{L} P_{L} + f_{2}^{R} P_{R} \right) t$$
 where, in the SM $f_{1}^{L} \approx 1$, $f_{2}^{L} = f_{1}^{R} = f_{2}^{R} = 0$ +h.c.



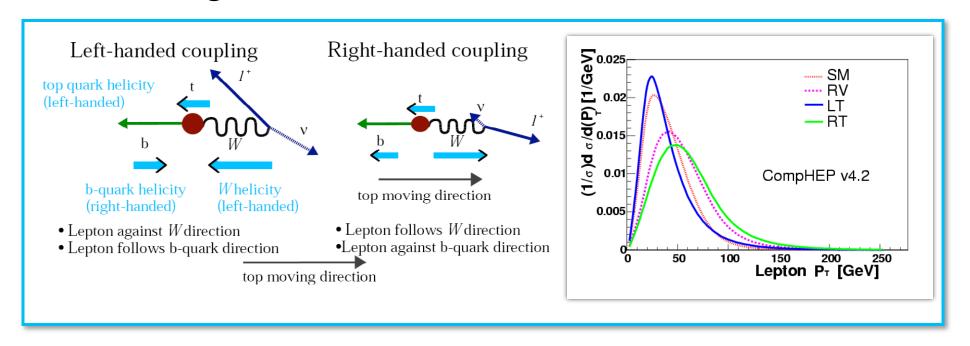
t-channel("tqb")

- Two non -zero couplings at a time
 - Consider 3 scenarios
 - Simultaneous limit on two couplings

Only
$$f_1L$$
, f_1R non-zero Only f_1L , f_2L non-zero Only f_1L , f_2R non-zero

anomalous couplings vs SM

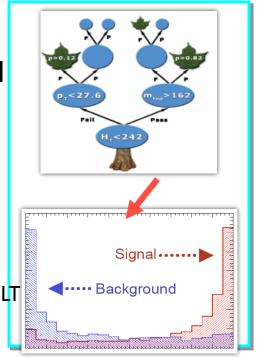
 presence of anomalous couplings changes the production cross-section, and kinematics and angular distributions



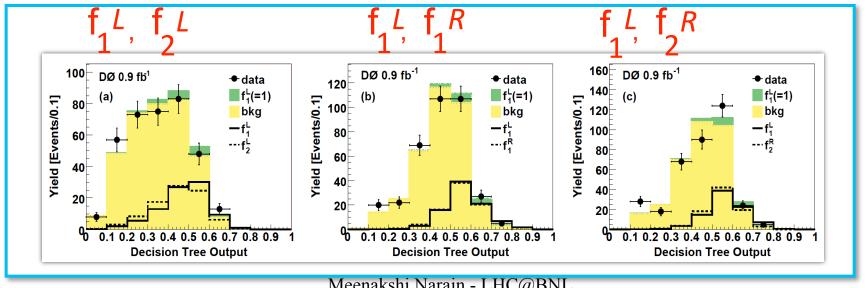
multivariate analysis

- Use Boosted Decision Trees to discriminate signal from background
- For every analysis, train 2 signals against sum of backgrounds

scenario: (tb + tqb)LV + (tb + tqb)RV scenario: (tb + tqb)LV + (tb + tqb)LT + (tb + tqb)LV+LT scenario: (tb + tqb)LV + (tb + tqb)RT



 $f = 1 \text{ fb}^{-1}$



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Limit Setting

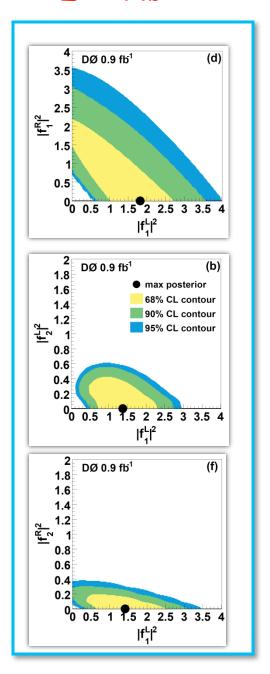
- Bayesian approach for limit setting
- Simultaneous limit setting for two signals by calculating 2 dimensional posterior probability density

Scenario	Cross Section	Coupling
(L_1,L_2)	$4.4^{+2.3}_{-2.5}$ pb	$ f_1^L ^2 = 1.4^{+0.6}_{-0.5}$ $ f_2^L ^2 < 0.5 \text{ at } 95\% \text{ C.L.}$
(L_1,R_1)	$5.2^{+2.6}_{-3.5}$ pb	$ f_1^L ^2 = 1.8^{+1.0}_{-1.3}$
(L_1,R_2)	$4.5^{+2.2}_{-2.2}$ pb	$ f_1^R ^2 < 2.5 \text{ at } 95\% \text{ C.L.}$ $ f_1^L ^2 = 1.4^{+0.9}_{-0.8}$
		$ f_2^R ^2 < 0.3$ at 95% C.L.

First experimental limits on tensor couplings!

(PRL 101, 221801 (2008))

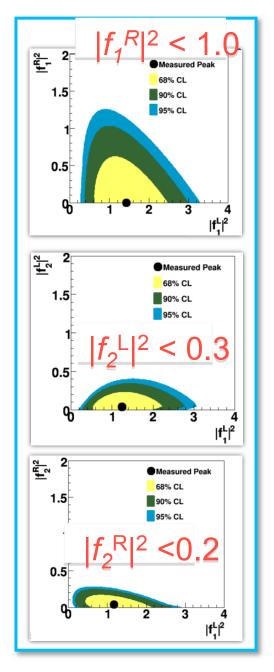
$\mathcal{L} = 1 \text{ fb}^{-1}$



Combination – tying it all together!

- W helicity measurement in top pair decays
- Anomalous couplings measurement in single top
- Bayesian analysis:
 - output of W helicity analysis forms input prior to single top anomalous couplings
- Observed posterior from data: single top and W helicity combined

 $\mathcal{L} = 1 \text{ fb}^{-1}$



analysis trends

- Looking at the whole picture
 - Tie together results in different channels and for different observables to compare the whole picture with predictions.
 - Mass vs cross section
 - Cross sections in different channels (t->Wb branching fraction, missing decay modes, charged Higgs, etc)
 - Wtb couplings (W helicity in ttbar decays and single top production/decay)

physics implications

- top quark is a known factor
 - Agrees with standard model (at least to precision probed at Tevatron – precision needed to estimate top as a background)
 - Can be simulated reliably with existing MC generators to estimate backgrounds to new physics
 - Provides important calibration point
 - b-tagging performance
 - jet energy scale
 - Tests of complex analysis chain of a know signal

physics opportunities

- top as a probe for new physics
 - New production modes
 - Top as a decay product
- measurements that are statistically limited at Tevatron
 - Rare top decays
 - Spin correlations
 - Single top production

conclusion

- top physics has come a long way since 1995
- top quark mass measured to 0.75%
 - reaching uncertainties below 1 GeV
- measurement of top properties and possible non-standard physics in t-W-b couplings are consistent with SM
- searches for new physics are conitnuing
- soon the torch will be passing to the LHC.

```
http://www-d0.fnal.gov/Run2Physics/top/
http://www-cdf.fnal.gov/physics/new/top/top.html
```

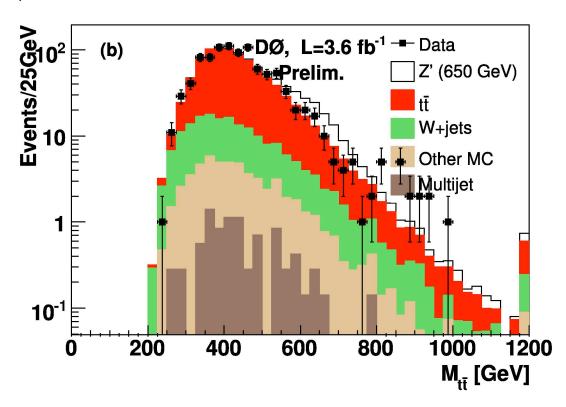


tt resonances

- D0 (3.6 fb⁻¹)
 - technicolor Z'→tt

Hill & Parke, PRD 49 (1994) 4454

 $M_{Z'}>820 \text{ GeV}$ for $\Gamma_{Z'}/M_{Z'}=1.2\%$



xsec syst

D0 I+jets b-tag

uncertainty source 0.15 pb vertex 0.11 pb e id μ id 0.08 pb 0.12 pb jet id non-W bkg 0.06 pb jet response 0.30 pb MC model 0.29 pb 0.48 pb b-tagging efficiency 0.69 pb total

kinematic likelihood

source	selection	fit	total
vertex	0.13 pb		0.13 pb
e id	0.10 pb		0.10 pb
μ id	0.06 pb		0.06 pb
jet id	0.10 pb	0.02 pb	0.12 pb
non-W bkg	0.10 pb		0.10 pb
jet response	0.35 pb	0.26 pb	0.11 pb
MC model	0.13 pb	0.09 pb	0.11 pb
template stats		0.15 pb	0.15 pb
total			0.36 pb

searches for non-standard physics

- quarks with charge 4/3e→ disfavored
- FB ttbar asymmetry → consistent with sm
- 4th generation t' quarks → m > 284 GeV
- scalar top production → no evidence
- charged Higgs bosons → limits on H⁺
- tb resonances →tb,t→H⁺b
- ttbar resonances
- FCNC decays of top quarks